

## Searches for charged Higgs bosons in ATLAS

P. CZODROWSKI on behalf of the ATLAS COLLABORATION

*Technische Universität Dresden - Dresden, Germany*

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**Summary.** — The presented analyses were performed by the Higgs to Complex States Working Group of the ATLAS Collaboration, using  $pp$  collisions data taken by the ATLAS detector in 2010 and 2011 at the LHC. An outline of the analysis and the result of the charged Higgs in the  $c\bar{s}$  final state, which uses kinematic fitting algorithms, was presented. Upper limits on  $\mathcal{B}(t \rightarrow bH^+) \times \mathcal{B}(H^+ \rightarrow c\bar{s}) = 0.25$  to  $0.14$  for  $m_{H^\pm} = 90$  to  $130$  GeV have been derived. In the second part the charged Higgs in the hadronic  $\tau$  final state analysis, which utilizes different techniques to estimate all background contributions from the data, was summarized. Values of the product of branching ratios  $\mathcal{B}(t \rightarrow bH^+) \times \mathcal{B}(H^+ \rightarrow \tau + \nu)$ , larger than  $0.03$  to  $0.10$  have been excluded in the  $H^\pm$  mass range  $90$  to  $160$  GeV.

PACS 14.65.Ha – Top quarks.

### 1. – Introduction

The charged Higgs boson is predicted by many extended Higgs scenarios, such as models containing Higgs triplets and Two-Higgs-Doublet Models (2HDM). The observation of charged Higgs bosons,  $H^\pm$ , would indicate physics beyond the Standard Model (SM). Both analyses presented here [1] consider the type-II 2HDM, which is also the Higgs sector of the Minimal Supersymmetric Standard Model (MSSM). For charged Higgs boson masses,  $m_{H^\pm}$ , smaller than the top quark mass,  $m_t$ , the dominant production mode at the LHC for  $H^\pm$  is through top quark decay via  $t \rightarrow bH^+$ . The dominant source of top quarks at the LHC is  $t\bar{t}$  production (fig. 1).

### 2. – The $c\bar{s}$ + jets final state

The goal of this analysis is to search for  $H^+ \rightarrow c\bar{s}$  decays. In the absence of a signal signature in data, exclusion limits for the top decay branching ratio to a light charged Higgs boson  $\mathcal{B}(t \rightarrow bH^+) \times \mathcal{B}(H^+ \rightarrow c\bar{s})$  are derived from early ATLAS data. A kinematic fit improves the sensitivity for a possible signal and allows to set very good limits, using only  $35 \text{ pb}^{-1}$ .

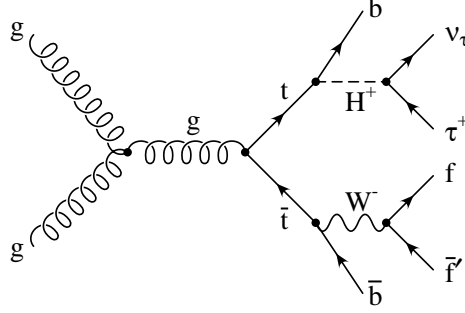


Fig. 1. – Example for a leading-order Feynman diagram for the production of a charged Higgs boson through gluon fusion in  $t\bar{t}$  decays.

**2.1. The kinematic fit.** – In lepton plus four jets events, the two jets originating from decays of  $H^+$  need to be identified in order to reconstruct the mass of  $H^+$  candidates. A kinematic fitter [2] is used to identify and reconstruct the mass of dijets from  $W/H^+$  candidates, by fully reconstructing the  $t\bar{t}$  system. In the kinematic fitter, the lepton, the missing transverse energy,  $E_T^{\text{miss}}$ , (from the neutrino), and four jets are assigned to the decay partons from the  $t\bar{t}$  system.

**2.2. Result and limits.** – The sensitivity of the analysis presented in [4] is comparable to the limits obtained at the Tevatron, where datasets with more than twenty-five times the integrated luminosity were used. The limits are shown in fig. 2 where the extracted limits are compared with the previous limits from the Tevatron [2, 3].

### 3. – The $\tau$ + jets final state

The goal of this analysis is to search for the light charged Higgs bosons in the decay channel  $H^+ \rightarrow \tau + \nu$ , all the backgrounds of the selection are estimated from data using data-driven techniques [5]. Only a selection of the methods applied is presented below. In absence of a significant signal, limits on the branching ratio of  $\mathcal{B}(t \rightarrow bH^+) \times \mathcal{B}(H^+ \rightarrow \tau + \nu)$  are derived.

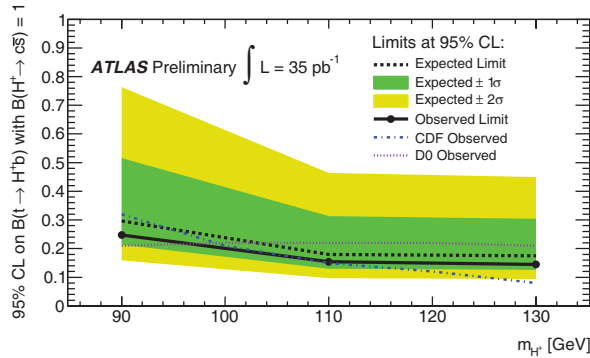


Fig. 2. – The extracted 95% CL upper limits on  $\mathcal{B}(t \rightarrow bH^+)$  from the ATLAS data are compared with the expected results and results from the Tevatron. These results assume  $\mathcal{B}(H^+ \rightarrow c\bar{s}) = 100\%$ . The ATLAS limits shown are calculated using the  $CL_s$  limit setting procedure.

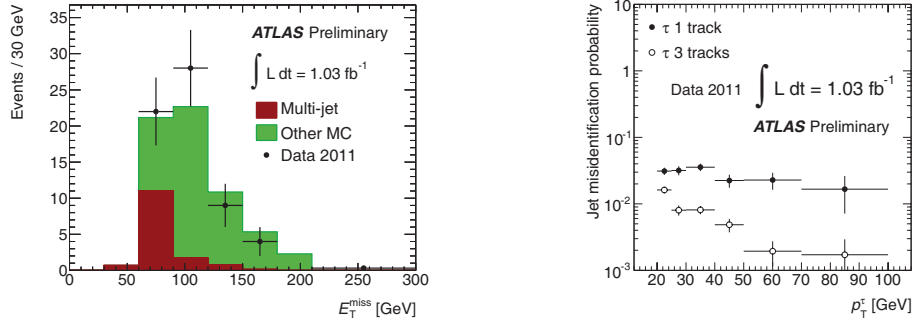


Fig. 3. – Left: multi-jet estimate: A fit to the  $E_T^{\text{miss}}$  distribution in data after all selection cuts using two shapes (one for the multi-jet model, and one for all other background processes, dominated by  $t\bar{t}$  and  $W$ +jets) is shown. The multi-jet fraction estimated after all selection cuts is  $(23 \pm 10)\%$ . Right: jet  $\rightarrow \tau$  misidentification probabilities measured from  $\gamma$ +jet events for 1-track and 3-track tight likelihood based  $\tau$  identification, *versus*  $p_T$ , using an integrated luminosity of  $1.03 \text{ fb}^{-1}$  are shown.

**3.1. Multi-jet background estimation and jet-to- $\tau$  misidentification probability.** – The multi-jet background is estimated by fitting its  $E_T^{\text{miss}}$  shape to data (fig. 3 (left)). In order to study this shape in a data-driven way, a control region is defined by inverting the  $\tau$  identification and  $b$ -tagging requirements. A measurement of the probability for a jet to be misidentified as a hadronically decaying  $\tau$  lepton is used to predict the yield of jet-to- $\tau$  misidentification events (fig. 3 (right)) from the most important SM backgrounds with intrinsic  $E_T^{\text{miss}}$ .

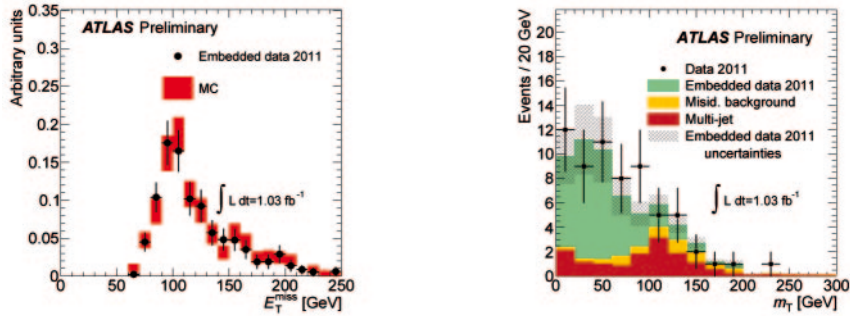


Fig. 4. – Left: The embedded data  $E_T^{\text{miss}}$  shape, normalized to unit area, is compared to  $t\bar{t}$  simulation after applying the  $H^+$  selection (shaded area indicates the statistical uncertainties on the MC simulation). Right: Comparison of the  $m_T$  shape for embedded events *versus* collision data. The prediction using the embedding method is stacked on top of the expected backgrounds with objects misidentified as  $\tau$  jets: MC expectation for  $t\bar{t}$  and electroweak processes, and the data-driven estimate for multi-jet events. The comparison is done after the  $H^+$  event selection and after normalizing the  $m_T$  distribution of embedded events to the data distribution in the range 0–40 GeV (gray area indicates the statistical and systematic uncertainties of the embedding method estimate).

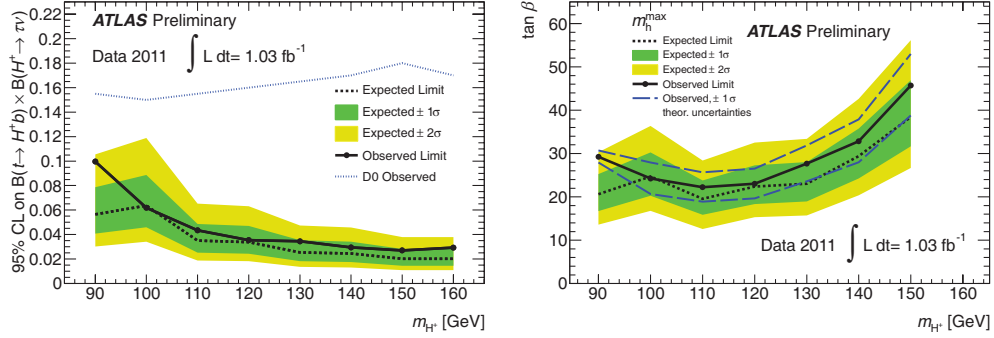


Fig. 5. – Left: Expected and observed 95% CL exclusion limits for charged Higgs boson production from top quark decays as a function of  $m_{H^+}$  in terms of  $\mathcal{B}(t \rightarrow H^+ b) \times \mathcal{B}(H^+ \rightarrow \tau^+ \nu)$ . For comparison, the best limit provided by the Tevatron experiments is shown [3]. Right: Limit for charged Higgs boson production from top quark decays in the  $m_{H^+}$ - $\tan \beta$  plane. Results are shown for the MSSM scenario  $m_h^{\max}$ .

**3.2. The embedding method.** – The method consists of collecting a control sample of real  $t\bar{t}$ , single-top and  $W$ +jets events with a muon and replacing the detector signature of this muon with that of a simulated  $\tau$  lepton. The reconstruction is re-applied to the new hybrid events which are then used to estimate the true- $\tau$  background of the selection (fig. 4). The whole event (except for the  $\tau$  jet) is taken directly from data, including the underlying event and pile-up, missing energy,  $b$  quark jets and light quark jets.

**3.3. Result and limits.** – Using data-driven background estimates, no statistically significant excess of events is observed in  $1.03 \text{ fb}^{-1}$  of collision data. Exclusion limits are set on the branching ratio  $t \rightarrow bH^+$  (fig. 5 (left)), and in the  $m_{H^+} - \tan \beta$  plane (fig. 5 (right)), by rejecting the signal hypothesis at the 95% confidence level applying the  $\text{CL}_S$  procedure [6, 7]. A profile likelihood ratio [8] is used with the  $m_T$  distribution as the discriminating variable. The final limits are based on the asymptotic distribution of the test statistic [8].

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